Cusp structures: combining multi-spacecraft observations with ground-based observations


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Abstract. Recent simultaneous observations of cusp structures with Polar, FAST and Interball revealed remarkably similar features at spacecraft crossing the cusp. Such stable cusp structures could be observed up to several hours only during stable solar wind conditions. Their similarities led to the conclusion that for such conditions large-scale cusp structures are spatial structures related to a global ionospheric convection pattern and not the result of temporal variations in reconnection parameters.

With the launch of the Cluster fleet we are now able to observe precipitating ion structures in the cusp with three spacecraft and identical instrumentation. The orbit configuration of the Cluster spacecraft allows for delay times between spacecraft of about 45 min in crossing the cusp. The compact configuration of three spacecraft at about the same altitude allows for the analysis of cusp structures in great detail and during changing solar wind conditions. Cluster observations on 25 July 2001 are combined with SuperDARN radar observations that are used to derive a convection pattern in the ionosphere. We found that large-scale cusp structures for this Cluster cusp crossing are in agreement with structures in the convection pattern and conclude that major cusp structures can be consistent with a spatial phenomenon.

Key words. Magnetospheric physics (energetic particles, precipitating, magnetopause, cusp arid and boundary layers; solar wind-magnetosphere interactions)

1 Introduction

Convincing evidence about magnetic reconnection between the interplanetary magnetic field (IMF) and the geomagnetic field has been accumulated with the observation of magnetosheath ions in the boundary layer inside the magnetopause (e.g. Paschmann et al., 1979; Sonnerup et al., 1981) and precipitating ions in the cusp (e.g. Reiff et al., 1977; Escoubet et al., 1997). The incoming magnetosheath distribution is truncated as it crosses the magnetopause and only a limited part of the initial magnetosheath distribution is transmitted across the magnetopause. This truncated distribution has a characteristic D-shaped distribution that has been predicted by Cowley (1982) and observed by, for example, Fuselier et al. (1991, 2000). Once reconnection occurred at the magnetopause, magnetosheath ions will stream continuously from the magnetosheath into the magnetosphere (e.g. Lockwood and Smith, 1993, 1994; Onsager et al., 1993). Because of the convection of the newly-opened magnetic field lines with the solar wind, spectra observed by a satellite in the cusp will have distinctive ion energy dispersions. Their appearance at an observing satellite will also depend on the satellite trajectory with respect to the convection direction. Rosenbauer et al. (1975) predicted that for reconnection of the geomagnetic field with a southward directed IMF, precipitating ions observed at low altitudes in the cusp should exhibit a velocity filter effect with lower energy ions convecting further poleward. Such a dispersion was indeed observed by Shelley et al. (1976). Newell and Meng (1995)
pointed out that such a decline of the average ion energy in the cusp with increasing latitude also includes additional effects. Plasma crossing the magnetopause close to the reconnection site experiences an acceleration which progressively declines as magnetic field lines convect to higher latitudes away from the merging point. As the magnetic field lines straighten out, they continue to convect poleward because the field lines are embedded in the solar wind flow, and ion acceleration decreases and turns into deceleration poleward of the cusp. In addition, with increasing latitude the magnetosheath velocity is increasingly directed away from the magnetosphere, causing fewer ions with progressively lower energy to make it down to the Earth’s ionosphere (Newell and Meng, 1991). Over the last two decades there has been a debate as to whether dayside reconnection is quasi-steady or transient (e.g. Newell and Sibeck, 1993; Lockwood et al., 1994 and the references therein). A smooth and continuous latitude dispersion in the cusp should be expected for a steady rate of reconnection at the magnetopause. However, satellite observations from this region show that the energy of precipitating ions is rarely smooth and continuous with increasing latitude but show complicated structures with variations in flux levels and sudden changes in the energy of the precipitating ions (e.g. Newell and Meng, 1991; Escoubet et al., 1992).

The existence of these steps in the ion energy dispersion, also known as “stepped” or “staircase” cusp ion signatures, has been predicted by Cowley et al. (1991) and Smith et al. (1992), based on a model by Cowley and Lockwood (1992) of how ionospheric convection is excited. In this pulsating cusp model (see also Lockwood and Smith, 1989; 1990), the cusp precipitation between the steps is the result of pulses of enhanced magnetopause reconnection. Steps are the result of changes in the reconnection rate at the magnetopause that creates neighboring flux tubes in the cusp with different time histories since reconnection (e.g. Lockwood and Smith, 1994). A significant characteristic of temporal steps is their convection with the open magnetic field lines under the joint action of magnetic tension and shocked solar wind flow. This creates an ever-changing structural profile of precipitating cusp ions.

Satellites crossing the boundary between a newly-opened flux tube and an older (but still open) one would encounter a step-down in the ion energy dispersion, while satellites crossing from an older flux tube into a nearby opened one would see a step-up ion energy signature. The type of cusp structure encounter by a satellite in the cusp also depends, therefore, on the satellite velocity relative to the convection velocity of the cusp structure. In general, fast moving low-altitude satellites (e.g. FAST) will overtake cusp structures, moving from a newer flux tube to an older one, and encounter step-down structures in the ion energy dispersion. In turn, a slow moving high-altitude satellite (e.g. Polar) will be overtaken by convecting cusp structures and, therefore, encounter step-up ion energy dispersions.

The pulsating cusp model was further supported by combining satellite observations with ground-based observations from the EISCAT radar (see Lockwood, 1995; Lockwood et al., 1995; Neudegg et al., 1999; Milan et al., 2000; McWilliams et al., 2001). Convecting flux tubes caused by reconnection pulses would move along the convection flow. Therefore, the observation of flow across a step in the cusp ion energy dispersion revealed the temporal nature of cusp structures. These cusp structures, the result of temporal variations of the reconnection rate at the magnetopause, are also discussed by Boudouridis et al. (2001). This model is based on the combination of the Bursty Single X-line Reconnection Model, together with the Multiple X-line Reconnection Model, to explain overlapping cusp steps observed by two DMSP spacecraft.

Flux tubes on open field lines with precipitating magnetosheath ions could also be spatially separated, emanating from multiple X-lines. Crossing the boundary into such a spatially separated different flux tube would also appear as a step in the ion energy dispersion, due to the different time history since reconnection for field lines within the two flux tubes (Lockwood et al., 1995). However, this step would not be convected with the solar wind but would appear as a standing feature in the cusp. Independent of the time delay between the cusp crossings or the satellite velocities, the satellites should encounter unchanged cusp structures at about the same latitude, observing a spatial feature. Such an observation would indicate that the reconnection rate at the magnetopause is rather stable and not highly variable, or even stops entirely for a limited period of time.

The appearance of spatial structures has been recently discussed by Wing et al. (2001), who modeled cusp precipitation characteristics for periods with a dominant IMF $B_y$ component. For these conditions they found that a characteristic “double cusp” signature was not only predicted but also observed in DMSP satellite data.

Also using DMSP observations, Onsager et al. (1995) showed two cusp crossings of the high-altitude Dynamic Explorer 1 (DE 1) and low-altitude DE 2 spacecraft separated by 20 min. A similar step in the ion dispersion signature at both spacecraft was interpreted as a spatial structure rather than a temporal variation of the reconnection rate. This event is especially interesting since the low-orbiting satellite encountered an upward step. A temporal convecting cusp structure would require the satellite to move along the open-closed field line boundary, to allow the convecting structure to overtake the low-altitude fast moving satellite. However, the observing satellite was in a meridional orbit.

To avoid the ambiguity of single satellite observations in distinguishing between spatial and temporal effects, Trattner et al. (1999, 2002a, b) have used pairs of Interball-Polar and FAST-Polar satellites to investigate the temporal or spatial nature of the cusp structure. They found that stepped ion distributions during stable solar wind and IMF conditions are not consistent with the pulsed reconnection model. Two magnetic conjunctions in the cusp by Interball and Polar revealed complicated cusp structures that appeared to be stable and unchanged for 1.5 h (Trattner et al., 1999). In a subsequent study, Trattner et al. (2002a) compared four cusp cross-
trations from three Cluster spacecraft with SuperDARN radar measurements (Lepping et al., 1995; Ogilvie et al., 1995). These data are provided by the ISTP key parameter web page.

The Cluster satellites carry the Cluster Ion Spectrometers (CIS) (Rème et al., 2001) that provide high-precision 3-D mass-resolving spectrometer capable of providing full 3-D distributions of the major ion species (H\(^+\), He\(^{2+}\), He\(^+\) and O\(^+\)) in the energy range from about 20 eV/e to 40 keV/e. CIS data are available on three spacecraft for this study.

The cusp observations presented in this study will only include H\(^+\) CIS/CODIF data from SC1, SC3 and SC4. We have selected a Cluster cusp crossing on 25 July 2001, from 23:00 UT to 26 July 2001, 00:25 UT, with relatively stable IMF conditions for the cusp crossing of SC1 and SC4, and with a significant change in IMF conditions for the cusp crossing of SC3. As in earlier studies (e.g. Trattner et al., 1999, Trattner et al., 2002a, b) we will investigate ion energy dispersions of multi-spacecraft cusp crossings and analyze the location and temporal changes of the cusp structures, i.e. major jumps in the cusp ion energy dispersion. Earlier observations with the FAST and Polar spacecraft (e.g. Trattner et al., 2002a) had the advantage of a large altitude separation between the observing satellites, resulting in significantly different satellite velocities while crossing the cusp and subsequently different observation times of structures inside the cusp. Spatial cusp structures would appear unchanged when observed by satellites at different altitudes and different velocities, but the nature of temporal structures would result in characteristic changes in appearance of cusp structures (step-up versus step-down). The Cluster satellites cross the cusp at about the same altitude that eliminates this advantage. However, this will be compensated for by using three identical instruments on board three satellites, crossing through the cusp with a time delay of up to 45 min and a spatial separation in MLT of about 50 min. The Cluster observations are supplemented with simultaneous radar observations by the 8 operating Northern Hemisphere SuperDARN radars which were running in a Cluster support mode for the selected cusp crossing.

In addition, Wind magnetic field and solar wind plasma data (MFI and SWE) are used as solar wind context measurements (Lepping et al., 1995; Ogilvie et al., 1995). These data are provided by the ISTP key parameter web page.
3 Observations

Figure 1 shows solar wind conditions for the Cluster cusp crossing on 25 July 2001, observed by the Wind solar wind analyzer (SWE) and the magnetic field investigation (MFI). The Wind spacecraft was located at about \( X_{\text{GSE}} = 45 R_E, Y_{\text{GSE}} = 264 R_E \) and \( Z_{\text{GSE}} = 18 R_E \). The solar wind data have been propagated by about 8 min to account for the travel time from the Wind spacecraft to the magnetopause. We have used the actual bow shock and the magnetopause locations in the travel time calculation by using the measured solar wind conditions with the Farris and Russell (1994) bow shock model and the Petrinec and Russell (1996) magnetopause model. The location of the Wind spacecraft (upstream, but far in the afternoon sector) is not ideal for the calculation of an accurate convection time of solar wind structures to the magnetopause. No ACE spacecraft observations were available to cross check the solar wind observations for this time interval. Figure 1 shows a solar wind density \( N \) of about 4 \( \text{cm}^{-3} \) (top panel) and a solar wind velocity \( V_x \) of about 560 km/s (middle panel). The IMF components \( B_x \) (black line), \( B_y \) (green line) and \( B_z \) (colored area) are shown in the bottom panel. At the beginning of the Cluster cusp crossing until about 23:40 UT, the IMF had a typical Parker spiral configuration with a positive \( B_x \) of about 3 nT, a negative \( B_y \) of about \(-4\) nT and a negative \( B_z \) of also about \(-4\) nT (blue colored area). At about 23:35 UT the \( B_z \) component decreased in strength to about 0 nT before increasing in strength and changing northward to about 4 nT (red colored area) at 23:47 UT. The \( B_y \) component changes direction from negative to positive for 15 min at about the same time as the \( B_z \) component changed to northward. The \( B_y \) component also switched from positive to negative for about 15 min during this time. Vertical black lines indicate the times when Cluster satellites crossed into the cusp to illustrate the temporal separation of the spacecraft.

Cluster satellites crossed into the cusp to illustrate the temporal separation of the spacecraft. SC4 and SC1 entered the cusp at 23:08 UT and 23:15 UT, respectively, during which the IMF was southward and stable for an extended period of time. SC3 entered the cusp 7 min after the IMF switched northward.

Plate 1 shows omnidirectional \( \text{H}^+ \) flux measurements (1/(cm\(^2\) s sr keV/e)) observed by the CIS instruments on three Cluster satellites for the cusp crossings on 25 July 2001. Fluxes from satellite SC1 (top panel), SC3 (middle panel) and SC4 (bottom panel), observed in an MLT range from 14:00 to 11:00, and invariant latitude (INVLAT) range from 76.8\(^\circ\) to 86\(^\circ\) and a geocentric distance from 4.8 \( R_E \) to 6 \( R_E \), are shown. White regions in the color-coded plot indicate regions with flux levels above the maximum indicated flux level in the color bars. SC1 enters the cusp at about 23:15 UT, marked by a sudden flux drop for the \( \text{H}^+ \) ions above 1 keV/e and a significant increase in ion energy and flux intensity below 1 keV/e (see white line 1a). SC1 subsequently observes the typical cusp ion energy dispersion for a southward interplanetary magnetic field, with lower energy ions arriving at higher latitudes (e.g. Reiff et al., 1977; Smith and Lockwood, 1996). The ion energy distribution decreases smoothly, indicating a constant magnetospheric reconnection rate at the magnetopause. At about 23:37 UT SC1 encounters a sudden increase in the ion energy dispersion (1c), consistent with a typical step-up ion signature that can occur if the satellite crossed onto magnetic field lines that were reconnected more recently. The ion energy of the precipitating ions again decreases until about 23:45 UT, where a new low is reached. Pitch-angle analysis of this low energy distribution shows that it is composed entirely of \( \text{H}^+ \) ion outflow onto field lines that form the plasma mantle.

SC3 crosses into the cusp at 23:54 UT, also indicated by a
Plate 1. Cluster/CIS observation for the cusp crossings on 25 July 2001. Plotted are H\(^+\) omnidirectional flux measurements (1/(cm\(^2\) sr s keV/e)) for satellite SC1, SC3 and SC4. All satellites encounter distinctive structures, sudden jumps in the ion energy dispersion that are similar on SC1 and SC4 but different on the later arriving SC3 satellite.

white line in the color spectrogram (3a). SC3 also observes a decreasing ion energy dispersion typical for a stable rate of reconnection with no further cusp structures later on.

The first Cluster satellite to enter the cusp on 25 July 2001, is SC4 at 23:08 UT (4a). The cusp encounter is also followed by a decreasing ion energy dispersion which is reversed at about 23:13 UT. The precipitating ion energy reaches a new maximum at 23:15 UT (4b), the same time as SC1 enters the cusp. The ion energy starts again to decrease before a sudden brief rise at 23:35 UT. A detailed pitch-angle analysis showed that this signature is caused by ionospheric H\(^+\) outflow (4d) and not ion precipitation from the magnetosheath. A pitch-angle analysis of the proton distribution for the same time interval revealed that such an localized ion outflow distribution was also present at the position of SC1 but not as clearly separated from the immediately following downward precipitating ions, as at the location of SC4. SC4 encounters this second sudden increase in ion energy about 23:37 UT (4c), similar to the increase observed by SC1 at about the same time. This step-up structure is also followed by a decrease in ion energy until about 23:45 UT, where a constant low-energy flux is reached, typical for high-latitude ion outflow.

Figure 2 shows the spatial separation of the Cluster spacecraft. Plotted are the magnetic footpoints of the Cluster satellites (SC1 blue, SC3 green, SC4 red) during their cusp crossing at 25 July 2001. Also shown are convection cells (black lines) derived from SuperDARN radar observations, averaged over the time interval from 23:00 UT to 23:30 UT. A larger separation of the contour lines indicates a slower flow velocity. Symbols along the magnetic footpoints of the Cluster satellites mark 15-min time segments. SC1 and SC4 are not only close together on the temporal scale (see Plate 1 and Fig. 1), entering the cusp within minutes of each other, but also have a small spatial separation, bringing their magnetic footpoints almost on top of each other. All Cluster magnetic footpoints are inside the dusk clockwise-rotating ionospheric convection cell (solid black lines) that covers the northern edge of North America, including Alaska and Greenland. The magnetic footpoints are also in the vicinity of an equatorward directed bulge in the dusk convection cell, where they finally cross into the dawn counterclockwise convection cell (dashed black lines) at high latitudes.

For orientation purposes, the white ellipsoids in Fig. 2 represent the average location of the auroral oval for the time of the Cluster cusp crossing on 25 July 2001. The location of the auroral oval is based on the model by Hardy et al. (1987) for \(K_p = 3\). The white dashed line in Fig. 2 represents the terminator. Both convection cells are located on the dayside. From the form of the clockwise convection cell we expect the location of the open-closed field line boundary to follow closely the location of the auroral oval.

Figure 3 presents color-coded line-of-sight velocity data (blue for flow towards, red for flow away) from the 8 operating Northern Hemisphere SuperDARN (Greenwald et al., 1995) radars, for the time interval 23:30–23:31 UT. The SuperDARN data coverage is especially good in the region of the equatorward directed bulge in the dusk convection cell, where the Cluster footpoints are located (compare with
Fig. 2. Magnetic footpoints of the Cluster satellites (SC1 blue, SC3 green, SC4 red) during their cusp crossing at 25 July 2001, to illustrate their spatial separation. Also shown are convection cells derived from SuperDARN radar observations for the time period from 23:00 UT to 23:30 UT. Each radar of the SuperDARN network was running a cluster support mode, comprising a 16 beam scan, with each beam sounding 70 range gates, each of 45 km, starting at a range of 180 km. The dwell time for each beam was 3 s, with scans synchronised to start at 1 min intervals. In addition to the line-of-sight data, the ionospheric equipotential flow streamlines have been calculated using the technique of Ruohoniemi and Baker (1998), and are presented as contour lines. Here the fit to the line-of-sight data is made to a 6th spherical harmonic expansion, with the fit stabilized by a statistical pattern keyed to the upstream IMF data from the Wind satellite, delayed by 8 min, to allow for the propagation time from the spacecraft to the magnetopause (Ruohoniemi and Greenwald, 1996). Such an equipotential map represents a best estimate of the global ionospheric flow response to the magnetopause processes sensed by the in situ spacecraft.

Figure 4 is a composite plot that combines the temporal and spatial separations of the Cluster spacecraft into one plot. The Cluster magnetic footpoints and the ionospheric convection stream lines for 25 July 2001 are shown, at 23:08 UT. Overlayed on the magnetic footpoints are 14-min wide sections of the Cluster/CIS flux measurements presented in Plate 1, which are centered on the actual position of the Cluster satellites at 23:08 UT. This representation shows the actual Cluster measurements in time at their proper spatial location where they have been observed. All satellites are in the dusk convection cell, with SC1 and SC4 close to an equatorward bulge of this cell. At 23:08 UT SC4 (red foot point line) crossed the ion open-closed field line boundary and entered the cusp where it encounters downward precipitating magnetosheath ions. The position of SC4 and the ion open-closed field line boundary is marked with a white triangle on SC4’s red magnetic foot point line. SC1, slightly delayed in space and time to SC4, was still on closed field lines as was SC3. Based on the center and form of the SuperDARN convection cell and the known location of the ion open-closed field line boundary, where SC4 crossed into the cusp, the most likely position of the ion open-closed field line boundary is indicated with a black dashed line. Note that the location of the ion open-closed field line boundary is based on the arrival of downward precipitating ions at Cluster. These ions need time to travel from the reconnection site to the observing satellite. The boundary is, therefore, located somewhat poleward of where such a boundary would have been placed from using electron precipitation or the examination of the SuperDARN data.

SuperDARN radars provide information on the open-closed field line boundary in two ways. The cusp region is known to be characterized by high spectral widths, and the boundary between high and low radar spectral width measurements can then act as a proxy for the open-closed field line boundary (e.g. Baker et al., 1990, 1995). In addition, the radar determination of the convection reversal boundary acts as a similar proxy. In this study the data coverage prevents a good determination of the spectral width boundary (although higher latitude data from Kodiak does show higher widths), but the convection reversal boundary can be determined from the map potential contours, essentially following a curve from the voltage minimum (for example, the region within the smallest closed contour shown by a solid line in...
Fig. 4. Composite plot of Cluster magnetic footpoints and ionospheric convection streamlines for 25 July 2001, at 23:08 UT. Overlayed on the magnetic footpoints are 14-min wide sections of the Cluster/CIS flux measurements presented in Plate 1, which are centered on the actual position of the Cluster satellites at 23:08 UT. Also indicated by a dashed line is the most likely position of the ion open-closed field line boundary based on the location of the Cluster SC4 satellite position, which entered the cusp at that time, and the SuperDARN convection pattern.

Such a boundary may be displaced from the actual open-closed field line boundary (for example, by viscous coupling effects) but does give the overall boundary shape and motion. Combined with the in situ spacecraft data from Cluster, however, the location and shape of the ion open-closed field line boundary may then be estimated.

In Fig. 4, the open-closed field line boundary determined from the convection cell method would closely follow the average location of the auroral oval (white line), while the “ion” open-closed field line boundary, depicted with a dashed black line and determined by the Cluster spacecraft crossing into the downward precipitating ion region, is located (as expected) poleward of this boundary.

Geomagnetic field lines reconnected at the magnetopause will start their convection cycle along the open-closed field line boundary. Following the convection lines evaluated from the SuperDARN radar data, a magnetic field line which starts a cycle at Point A will need to convect to B, to intersect the SC4 Cluster satellite. For this specific configuration, SC4 will travel only a short distance to higher latitudes (from the entry point into the cusp to point B) to intersect magnetic field lines that have been open for a rapidly increasing amount of time. Due to this rapidly increasing distance between the observing satellite and the ion open-closed field line boundary along the convection lines, this should result in a significant decrease in the energy of precipitating ions observed by SC4 with increasing invariant latitude. Plate 1 shows that this is indeed the case.

Figure 5 is a composite plot with the same format as Fig. 4, but for a later time. The Cluster magnetic footpoints and ionospheric convection streamlines for 25 July 2001 are plotted, at 23:15 UT. Overlayed on the magnetic footpoints are again 14-min wide sections of the Cluster/CIS flux measurements presented in Plate 1, centered on the actual position of the Cluster satellites at 23:15 UT. All Cluster satellites are still magnetically connected to the dusk ionospheric convection cell with SC4 inside the cusp, SC1 at the ion open-closed field line boundary and SC3 on closed field lines. The location of SC1 at the boundary is marked with a star along the blue magnetic foot point line that coincides with the position of the triangle where SC4 entered the cusp. Between the time when SC4 and SC1 crossed the ion open-closed field line boundary, the boundary did not move significantly. Based on this new position measurement of the equatorward position of the cusp and the center and form of the convection cell, the estimated position of the ion open-closed field line boundary is again marked with a dashed line. The equatorward bulge in the convection cell is less prominent at 23:15 UT compared to six minutes before when SC4 entered the cusp. This
change caused the convection streamlines to severely shorten between the ion open-closed field line boundary and the magnetic footpoints of the Cluster spacecraft. While position B along the magnetic footpoints of SC4 remained unchanged, position A at the boundary has moved considerable closer. Newly-opened magnetic field lines will not have to convect far to intersect the Cluster satellites, which, in turn, will allow precipitating ions with higher energies to be observed at the satellite. While SC1 has just entered the cusp, encountering precipitating ions injected onto the open magnetic field line with the highest energies, the initially decreasing ion energy dispersion profile for SC4 has turned around and increased in energy, forming a new maximum at 23:15 UT (4b in Plate 1 and Fig. 5). The new maximum in the ion energy dispersion profile observed at SC4 is in agreement with the observed shortening of the convection distance derived from the SuperDARN radar data.

Figure 6 shows the composite plot of Fig. 4 at 23:37 UT. SC1 and SC4 are deep inside the cusp, while SC3 is still on closed field lines. The original entry point of SC1 and SC4 into the cusp are marked with a star and a triangle along the tracks of their magnetic footpoints. At 23:37 UT the IMF shows a strong decrease in the value of $B_z$, which will subsequently result in a reversal from negative to positive $B_z$. The dusk convection cell has been elongated, lengthening the convection path again. The equatorward directed bulge has reasserted itself, moved rapidly equatorward, which, in turn, allowed for the dawn convection cell to move equatorward as well. At 23:37 UT SC1 and SC4 have progressed far enough poleward to be overtaken by the equatorward moving dawn convection cell. The transition from one convection cell to another resulted in an almost simultaneous sudden increase in the ion energy dispersion on both satellites, indicating that the ion open-closed field line boundary in the dawn convection cell is much closer to the SC1 and SC4 magnetic footpoints than in the dusk convection cell. The satellite positions at 23:37 UT, which are also the positions of the sudden increase in the ion energy dispersion, are marked with a small star (SC1) and a small triangle (SC4).

The sudden increase in the ion energy dispersion that coincides with a satellite moving into a neighbouring and spatially separated flux tube (or convection cell) was discussed by Trattner et al. (2002a, b) based on Polar and FAST precipitating ion observations of stable, unchanging cusp structures during stable solar wind conditions. Figure 6 shows not only that such a scenario can take place, but also that it occurs during dynamic solar wind IMF conditions. The change in
IMF conditions most likely caused a change in the location of the reconnection site and the associated convection pattern, which, in turn, caused a shift in the positions of the spatially separated flux tubes emanating from these reconnection sites.

Figure 7 shows the ionospheric convection streamlines and flux measurements for 25 July 2001, at 23:54 UT. SC1 and SC4 observe low energy ion precipitation, typical for plasma entering the open field lines at high magnetic latitudes poleward of the cusp. SC3 has finally reached the cusp and encounters the typical cusp ion energy dispersion, with higher energy ions arriving at subsequently higher latitudes. Using the position of the equatorward boundary of the cusp from SC3, and the center and form of the convection cell, we have estimated the location of the ion open-closed field line boundary. At 23:54 UT, the IMF has turned northward. The dawn convection cell was almost completely annihilated by this IMF change. The dusk convection cell is about to break up into smaller cells (not shown here), but at 23:54 UT the outlines of the original cell from the beginning of the Cluster cusp event are still visible. The estimated ion open-closed field line boundary also crosses the position of the boundary encountered by SC1 and SC4 almost 45 min earlier. These positions are marked with a white star and triangle symbols along their respective magnetic footpoints. Additional symbols along the magnetic footpoints of SC1 and SC4 at higher latitudes mark the position where the sudden step-up was encountered.

4 Summary and conclusion

Combining the observations of several satellites to separate spatial and temporal effects has proven to provide new insights into old ambiguities and is becoming increasingly more exploited since the launch of the Cluster spacecraft. Earlier studies were relying on chance conjunctions of often very different satellites in the cusps to study the temporal and spatial nature of cusp structures (e.g., Onsager et al., 1995; Trattner et al., 2002a). For cusp structures the use of satellites operating on vastly different altitudes is considered to be an advantage, since temporal features would appear differently compared to spatial features on the observing satellites. These earlier studies indeed revealed that, during stable solar wind conditions, structures in the cusp appear to be spatial features, most probably related to spatially separated flux tubes like, for example, different convection cells.

However, if the above conclusions are correct, then cusp structures should be spatial, even during changing solar wind and IMF conditions. The shift in the location of the reconnection line (or reconnection lines for multiple X-lines) at the magnetopause would certainly cause a change in the convection pattern in the ionosphere. New flux tubes will be opened or existing flux tubes moved to different locations. Even if these flux tubes are changing in time and location, cusp steps should be observed when satellites cross the boundaries between them. Therefore, we would have a temporally chang-...
wood and Smith, 1989, 1990, 1994; Lockwood et al., 1998). However, it does not rule out the existence of such cusp structures caused by the temporal variation of the reconnection rate. As mentioned above, the signature of such temporal reconnection pulses should show up as moving structures within one convection cell and would be in agreement with the observation of poleward moving transients in radar data.

In this study we have successfully linked satellite observations with an ionospheric convection pattern derived from SuperDARN radar observations. The combination of temporal and spatial scale lengths into one composite picture showed a remarkably good correlation between the encounter of cusp structures by the Cluster satellites and the motion of large-scale convection cells. The cusp structures observed during the 25 July 2001 Cluster cusp crossing are consistent with encountering variations of the location of spatially separated flux tubes.

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References


